# **Efficient Matching and Merging with Sector-Antenna Showers**

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### Outline

- 1) Antenna showers on sectorised phase spaces [Brooks, CTP, Skands 2003.00702]
- 2) Efficient (CKKW-L-style) merging with sector showers [Brooks, CTP 2008.09468]
  - (POWHEG also possible but not shown here; see [Höche, Mrenna, Payne, CTP, Skands 2106.10987])
- 3) Towards NNLO+PS matching with sector showers [Campbell, Höche, Li, CTP, Skands 2108.07133]

### VINCIA overview

# C. T. Preuss, **P. Skands**, R. Verheyen



- ullet originally developed as plug-in to Pythia 8.2 (started  $\sim$  2007 by P. Skands)
- $\bullet$  now part of  $\ensuremath{\mathrm{PYTHIA}}$  8.3 (since October 2019) as one of three showers:
  - "simple" shower (Pythia's  $p_{\perp}$ -ordered DGLAP shower)
  - VINCIA
  - Dire
- full-fledged antenna shower for ISR, FSR, coloured resonances (top)
- exact treatment of mass corrections (phase space and antenna functions)
- full helicity dependence in shower and MECs
- dedicated default tuning (similar to Pythia's Monash tune)

### Recent and ongoing developments:

- interleaved resonance decays for top, Z, W (P. Skands, R. Verheyen)
- interleaved coherent QED multipole shower (P. Skands, R. Verheyen)
- full-fledged (collinear) EW shower module (P. Skands, R. Verheyen)
- ullet QCD sector showers o this talk
- ullet efficient merging (CKKW-L) with sector showers o this talk
- ullet automated matrix element corrections (MECs) o this talk
- ullet (N)NLO matching with sector showers o this talk

# What is an event generator?

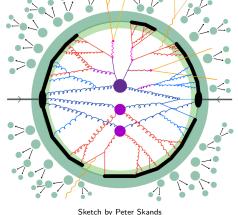
Particle-level event generators aim at simulating high-energy particle collisions in full detail by dividing events into three energy regimes:

Hard regime (multiple) high-energy  $2 \rightarrow n$  processes with small n

### Soft regime

forming and fragmentation of (visible) hadrons at low energies

### Transition regime QCD bremsstrahlung (+ QED/EW emissions)



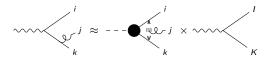


The "big players": PYTHIA, SHERPA, HERWIG

### Parton showers

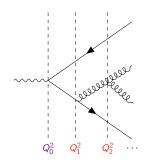
Parton showers dress a LO calculation with additional radiation, describing the evolution from parton level (quarks, gluons, ...) to the particle level (hadrons).

 amplitudes factorise in limits where emissions are soft  $(E_i \rightarrow 0)$  or collinear  $(\vartheta_{ik} \rightarrow 0)$ 



• starting from a **high scale**  $Q_0^2$ , further radiation is modelled under the assumption that it is soft/collinear and ordered

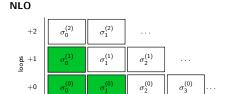
$$(Q_0^2 > Q_1^2 > Q_2^2 > \dots > Q_{\rm had}^2)$$



ullet evolves event from hard scale  ${\cal Q}_0^2$  to soft scale  ${\cal Q}_{
m had}^2$  but introduces **logarithms** of the form

$$lpha_{
m S}^{\it n} 
ightarrow lpha_{
m S}^{\it n} \log^{\it m} \left( rac{Q_0^2}{{
m Q}^2} 
ight) \, , {\it n} \leq 2{\it m} \, , \quad {
m large} \, {
m if} \, \, {
m Q}^2 \ll {\it Q}_0^2$$

### Parton showers vs fixed-order calculations



+1

### 

 $\sigma_1^{(0)}$ 

+1

 $\sigma_2^{(0)}$ 

+2

 $\sigma_0^{(0)}$ 

+0

+0

### Fixed-order calculations $\rightarrow$ hard jets

+0

- reliable at high scales if no large scale hierarchies are present
- accurate predictions for limited number of legs (+ loops)

+2

+3

• determines perturbative accuracy (LO, NLO, NNLO, ...)

### $\textcolor{red}{\textbf{Showers}} \rightarrow \mathsf{jet} \ \mathsf{substructure}$

- reliable in soft/collinear regions if large scale hierarchies are present
- approximate predictions for many particles
- determines logarithmic accuracy (LL, NLL, NNLL, ...)
- ⇒ largely complementary, so ideally combine them!



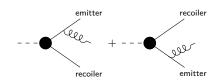
 $\sigma_3^{(0)}$ 

+3



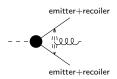
Many ways to skin a cat...

### **Dipoles**



- e.g. Sherpa CSS, Herwig dipole, Dire
- distinguish emitter and recoiler
- two branching kernels per colour dipole
- partition soft eikonal
- related to NLO dipole subtraction [Catani, Seymour hep-ph/9605323]

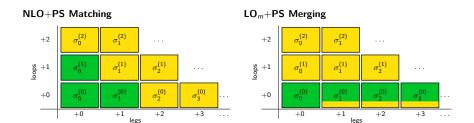
### Antennae



- e.g. Ariadne, Vincia
- both parents absorb transverse recoil
- one branching kernel per colour dipole
- partition collinear singularities
- related to (N)NLO antenna subtraction [Campbell, Cullen, Glover hep-ph/9809429]
   [Gehrmann-de Ridder, Gehrmann, Glover hep-ph/0505111]

Not mentioned here: **DGLAP** showers (Pythia  $p_{\perp}$ , Herwig  $\tilde{q}$ )

# Combining showers and fixed-order calculations



### Some disambiguation:

Matching combine a fixed-order (typically NLO) calculation with a parton shower, avoiding double-counting in overlap regions

Merging combine multiple inclusive (N)LO event samples into a single inclusive one with additional shower radiation, accounting for Sudakov suppression and avoiding double-counting in overlap regions (typically via phase-space slicing)

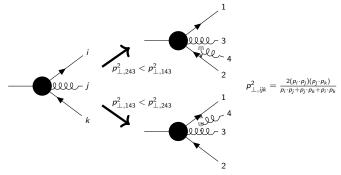
⇒ can we define a shower **designed** for **matching and merging**?

Part I: Sector-Antenna Showers

# Sector showers [Brooks, CTP, Skands 2003.00702; Lopez-Villarejo, Skands 1109.3608]

Idea: combine antenna shower with deterministic jet-clustering algorithm

• let shower only generate emissions that would be clustered by a  $(3 \mapsto 2)$  jet algorithm ( $\sim$  Arclus [Lönnblad Z.Phys.C 58 (1993)])



- ⇒ softest gluon always regarded as the emitted one
- ⇒ only one (most singular) splitting kernel contributes per phase space point

Since Pythia 8.304: full-fledged\* implementation of sector showers in Vincia

<sup>\*</sup>including FSR, ISR, resonance-decay showers

# Interlude: why sector showers?

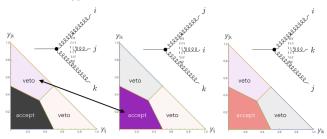
A variety of showers has been developed to date, so why another one?

- conceptually simpler structure (only a single branching per PS point)
- clear scale definitions due to deterministic clusterings
- simple removal of **overlaps**
- ⇒ full numerical control over shower domain & evolution
- ⇒ sector showers are **designed** for matching & merging

# Phase space sectors

Branching phase space gets divided into non-overlapping sectors.

• e.g. first emission in  $H \rightarrow gg$ :



- branchings in the shower are accepted if and only if they correspond to the correct sector
- $\bullet$  sectors defined by minimal  $p_{\perp}$  in event, but always contain:
  - ▶ soft endpoint
  - "full" collinear region for qg
  - "half" of the collinear region for gg with boundary at  $z = \frac{1}{2}$

Note: in general, non-trivial sector boundaries away from the singular limits!

### Sector antenna functions

Splitting kernels have to incorporate full single-unresolved limits for given PS point (KOSOWER subtraction terms [Kosower PRD 57 (1998) 5410, PRD 71 (2005) 045016])

• e.g. (FF)  $qg \mapsto qgg (s_{ij} = 2p_i \cdot p_j)$ :

$$A_{qg\mapsto qgg}^{\mathrm{sct}}(i_q, j_g, k_g) \rightarrow \begin{cases} \frac{2s_{jk}}{s_{ij}s_{jk}} & \text{if } j_g \text{ soft} \\ \frac{1}{s_{ij}} \frac{1+z^2}{1-z} & \text{if } i_q \parallel j_g \\ \frac{1}{s_{jk}} \frac{2(1-z(1-z))^2}{z(1-z)} & \text{if } j_g \parallel k_g \end{cases}$$

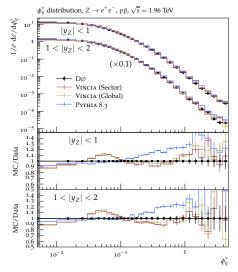
Compare to **global** antenna functions (a.k.a. *sub-antenna* functions):

• only "half" of the  $j_g \parallel k_g$  limit contained in the splitting kernel:

$$A_{qg \mapsto qgg}^{\mathrm{gl}}(i_q, j_g, k_g) \rightarrow \begin{cases} \frac{2s_{ik}}{s_{ij}s_{jk}} & \text{if } j_g \text{ soft} \\ \frac{1}{s_{ij}} \frac{1+z^2}{1-z} & \text{if } i_q \parallel j_g \\ \frac{1}{s_{ik}} \frac{1+z^3}{1-z} & \text{if } j_g \parallel k_g \end{cases}$$

ullet "rest" of the jk-collinear limit reproduced by neighbouring antenna  $(z \leftrightarrow 1-z)$ 

# Sector showers vs global showers

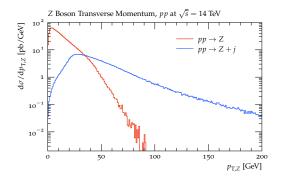


The sector approach is merely an **alternative way** to fraction singularities, so **formal accuracy** of the shower should be **retained**.

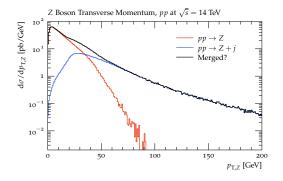
Note: same "global shower" tune in  $\mathrm{V}{\scriptstyle\mathrm{INCIA}},$  no MECs here

Part II: Efficient Merging with Sector Showers

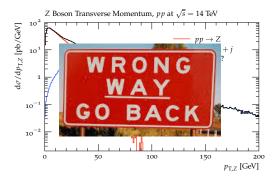
**Merging:** introduce (arbitrary) **merging scale** and let each calculation populate the phase space where it does best:



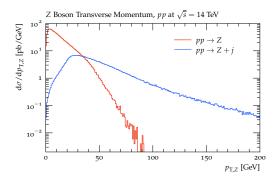
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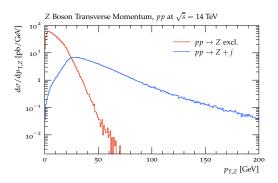
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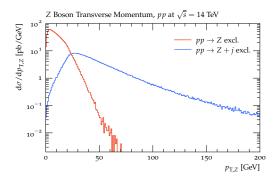
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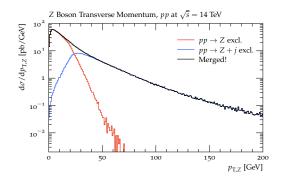
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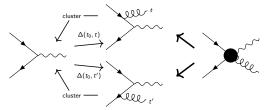
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# Merging with traditional showers: CKKW-L

Basic CKKW-L idea [Catani, Krauss, Kuhn, Webber hep-ph/0109231], [Lönnblad hep-ph/0112284]

- construct all possible shower histories, choose most likely
- let (truncated) trial showers generate Sudakov factors
- re-weight event by Sudakov factors



• number of histories scales factorially with number of legs

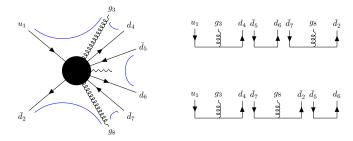
Number of Histories for n Branchings							
	n = 1	n = 2	n = 3	n=4	n = 5	n = 6	n = 7
CS Dipole	2	8	48	384	3840	46080	645120
Global Antenna	1	2	6	24	120	720	5040

quickly increasing complexity with multiplicity!

# Merging with sector showers (MESS) [Brooks, CTP 2008.09468]

Tree-level merging with sector showers straight-forward: start from CKKW-L and modify history construction (could be extended to NLO)

- sector showers have a single (!) history for gluon emissions at LC
- ullet to account for **gluon splittings**  $g\mapsto qar q$ , find all viable quark permutations

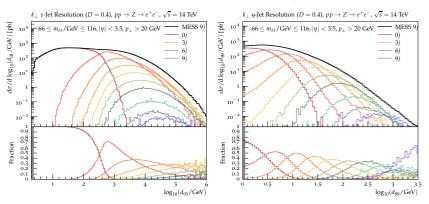


- for each colour-ordering, shower history again uniquely defined by sectors
- if multiple colour-orderings possible, choose one that maximises branching probability

Since Pythia 8.304: sector merging available with Vincia

# Merging with sector showers: validation

**Parton-level** results for merging in  $pp \rightarrow Z$  with up to **9 jets** (using HDF5 event samples from [Höche, Prestel, Schulz 1905.05120])



⇒ smooth transitions, no "sector effects" visible

# Merging with sector showers: efficiency

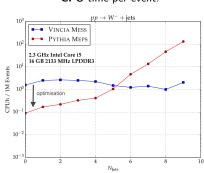
Gauge efficiency gains in  $pp \rightarrow Z+9j$  merging @ parton level (using HDF5 event samples from [Höche, Prestel, Schulz 1905.05120]).

# **memory** allocation/deallocation: Exclusive Contributions to $pp \rightarrow Z + 10$ jets

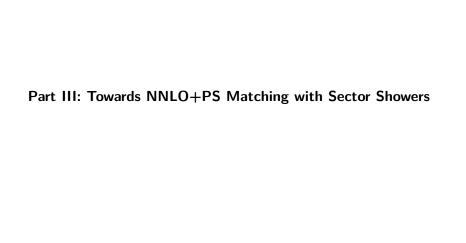
# BD 100 PATHIA MESS PATHIA MESS

Number of lets N<sub>tot</sub>

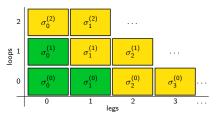
### CPU time per event:



- ⇒ ~ constant runtime and memory footprint in multi-jet merging
- ⇒ overall optimisation of the sector shower possible



# NLO+PS matching



Strategy developed  $\gtrsim 20$  years ago [Norrbin, Sjöstrand hep-ph/0010012] nowadays known as POWHEG matching [Nason hep-ph/0409146]

Alternative strategy: MC@NLO [Frixione, Webber hep-ph/0204244] (not discussed here)

• POWHEG master formula (for 2 Born jets):

• main trick: matrix-element correction (MEC) in first shower emission

$$\mathcal{S}_2(t_0, O) = \Delta_2(t_0, t_{
m c}) \mathcal{O}(\Phi_2) + \int\limits_t^{t_0} {
m d} \Phi_{+1} \, rac{{
m R}(\Phi_2, \Phi_{+1})}{{
m B}(\Phi_2)} \Delta_2(t, t_{
m c}) \mathcal{O}(\Phi_2)$$

where

$$\Delta_2(t,t') = \exp\left(-\int_{t'}^t \mathsf{d}\Phi_{+1} \, A_{2\mapsto 3}(\Phi_{+1}) w_{2\mapsto 3}^{\mathrm{MEC}}(\Phi_2,\Phi_{+1})\right) \, , \, \, w_{2\mapsto 3}^{\mathrm{MEC}} = \frac{\mathrm{R}(\Phi_2,\Phi_{+1})}{A_{2\mapsto 3}(\Phi_{+1})\mathrm{B}(\Phi_2)}$$

# Towards NNLO+PS [Campbell, Höche, Li, CTP, Skands 2108.07133]

**Idea**: "POWHEG at NNLO" (focus here on  $e^+e^- \rightarrow 2j$ )

$$\langle O \rangle_{\mathrm{NNLO+PS}} = \int \mathsf{d}\Phi_2 \, \mathrm{B}(\Phi_2) \left[ k_{\mathrm{NNLO}}(\Phi_2) \right] \left[ \mathcal{S}_2(t_0, O) \right]_{\mathrm{shower operator}}$$

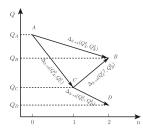
Key aspect: 2-particle and 3-particle Sudakovs with (N)LO MECs in  $\mathcal{S}_2$ 

$$\begin{split} \Delta_2^{\mathrm{NLO}}(t_0,t) &= \exp\bigg\{ - \int_t^{t_0} \mathsf{d}\Phi_{+1} \, \mathrm{A}_{2\mapsto 3}^{(0)}(\Phi_{+1}) w_{2\mapsto 3}^{\mathrm{NLO}}(\Phi_2,\Phi_{+1}) \bigg\} \\ &\times \exp\bigg\{ - \int_t^{t_0} \mathsf{d}\Phi_{+2}^{>} \, \mathrm{A}_{2\mapsto 4}^{(0)}(\Phi_{+2}) w_{2\mapsto 4}^{\mathrm{LO}}(\Phi_2,\Phi_{+2}) \bigg\} \\ \Delta_3^{\mathrm{LO}}(t,t') &= \exp\bigg\{ - \int_{t'}^t \mathsf{d}\Phi_{+1}' \, \mathrm{A}_{3\mapsto 4}^{(0)}(\Phi_{+1}') w_{3\mapsto 4}^{\mathrm{LO}}(\Phi_3,\Phi_{+1}') \bigg\} \end{split}$$

Divide double-emission phase space into **strongly-ordered** and **unordered** region:

$$d\Phi_{+2} = \underbrace{d\Phi_{+2}^{>}}_{u.o.} + \underbrace{d\Phi_{+2}^{<}}_{s.o.}$$

s.o. region: only **single-unresolved** limits u.o. region: only **double-unresolved** limits



### Towards NNLO+PS: MECs

Iterated tree-level MECs in ordered region [Giele, Kosower, Skands 1102.2126], [Fischer, Prestel 1706.06218]:

$$\begin{split} w_{2\mapsto3}^{\mathrm{LO}}(\Phi_2,\Phi_{+1}) &= \frac{\mathrm{R}(\Phi_2,\Phi_{+1})}{\mathrm{A}_{2\mapsto3}^{(0)}(\Phi_{+1})\mathrm{B}(\Phi_2)} \\ w_{3\mapsto4}^{\mathrm{LO}}(\Phi_3,\Phi_{+1}') &= \frac{\mathrm{RR}(\Phi_3,\Phi_{+1}')}{\mathrm{A}_{3\mapsto4}^{(0)}(\Phi_{+1}')\mathrm{R}(\Phi_3)} \end{split}$$

Tree-level MECs in unordered region:

$$w_{2\mapsto 4}^{\text{LO}}(\Phi_2, \Phi_{+2}) = \frac{\text{RR}(\Phi_2, \Phi_{+2})}{A_{2\mapsto 4}^{(0)}(\Phi_{+2})B(\Phi_2)}$$

**NLO** MECs for +1j state [Hartgring, Laenen, Skands 1303.4974]:

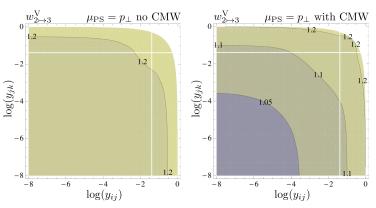
$$\begin{split} & w_{2 \mapsto 3}^{\mathrm{NLO}}(\Phi_{2}, \Phi_{+1}) = w_{2 \mapsto 3}^{\mathrm{LO}}(\Phi_{2}, \Phi_{+1}) \times (1 + w_{2 \mapsto 3}^{\mathrm{V}}(\Phi_{2}, \Phi_{+1})) \\ & w_{2 \mapsto 3}^{\mathrm{V}}(\Phi_{2}, \Phi_{+1}) = \left(\frac{\mathrm{RV}(\Phi_{2}, \Phi_{+1})}{\mathrm{R}(\Phi_{2}, \Phi_{+1})} + \frac{\mathrm{I}^{\mathrm{NLO}}(\Phi_{2}, \Phi_{+1})}{\mathrm{R}(\Phi_{2}, \Phi_{+1})} \right. \\ & \mathrm{NLO~Born} + 1j \qquad + \int_{0}^{t} \mathrm{d}\Phi'_{+1} \left[\frac{\mathrm{RR}(\Phi_{2}, \Phi_{+1}, \Phi'_{+1})}{\mathrm{R}(\Phi_{2}, \Phi_{+1})} - \frac{\mathrm{S}^{\mathrm{NLO}}(\Phi_{2}, \Phi_{+1}, \Phi'_{+1})}{\mathrm{R}(\Phi_{2}, \Phi_{+1})}\right] \right) \\ & \mathrm{NLO~Born} \quad - \left(\frac{\mathrm{V}(\Phi_{2})}{\mathrm{B}(\Phi_{2})} + \frac{\mathrm{I}^{\mathrm{NLO}}(\Phi_{2})}{\mathrm{B}(\Phi_{2})} + \int_{0}^{t_{0}} \mathrm{d}\Phi'_{+1} \left[\frac{\mathrm{R}(\Phi_{2}, \Phi'_{+1}, \Phi'_{+1})}{\mathrm{B}(\Phi_{2})} - \frac{\mathrm{S}^{\mathrm{NLO}}(\Phi_{2}, \Phi'_{+1})}{\mathrm{B}(\Phi_{2})}\right] \right) \\ & \mathrm{shower} \quad + \left(\frac{\alpha_{\mathrm{S}}}{2\pi} \log \left(\frac{\kappa^{2}\mu_{\mathrm{PS}}^{2}}{\mu_{\mathrm{T}}^{2}}\right) + \int_{0}^{t_{0}} \mathrm{d}\Phi'_{+1} \, \mathrm{A}_{2 \mapsto 3}^{(0)}(\Phi'_{+1}) w_{2 \mapsto 3}^{\mathrm{LO}}(\Phi_{2}, \Phi'_{+1}) \right) \end{split}$$

### Towards NNLO+PS: real-virtual corrections

Real-virtual correction factor

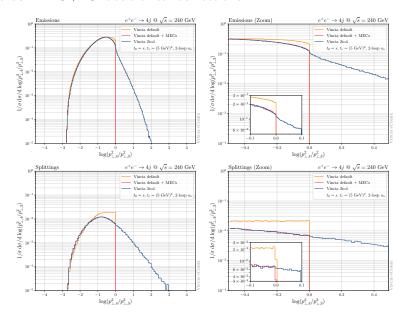
$$w_{2\mapsto3}^{\rm NLO}=w_{2\mapsto3}^{\rm LO}\left(1+w_{2\mapsto3}^{\rm V}\right)$$

studied in detail for  $Z o q \bar{q}$  in [Hartgring, Laenen, Skands 1303.4974]:



 $\Rightarrow$  now: generalisation & (semi-)automation in VINCIA in progress (using run-time interfaces to MCFM and Sherpa/Comix)

### Towards NNLO+PS: double-real corrections



### Conclusions

### Sector showers combine shower evolution with jet clustering to become maximally bijective

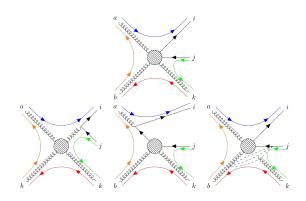
- "sectorised" VINCIA well validated against "global" VINCIA and PYTHIA (discontinuities? still searching...)
- ullet sector merging has  $\sim\!$  constant overall run time and memory usage
- sector showers default option in VINCIA as of PYTHIA 8.304

### This is just the beginning...

- sector merging easily extendable to NLO (lack of time that it hasn't been done yet...)
- sector decomposition facilitates inclusion of NLO antenna functions in shower evolution (including direct 2 → 4 branchings covering double-unresolved limits)
- antenna-based (N)NLO matching and shower evolution at NLO ongoing developments (currently on a proof-of-concept level for  $e^+e^- \rightarrow 2j$ , but can be extended!)

# **Backup**

### Sector definitions



For massless particles, the sector resolution is defined by:

$$Q_{\mathsf{res},j}^2 = \begin{cases} \frac{s_{ij}s_{jk}}{s_{ijk}} & \text{if } j \text{ is a } g \\ s_{ij}\sqrt{\frac{s_{jk}}{s_{ijk}}} & \text{if } (i,j) \text{ is a } q\bar{q} \text{ pair} \end{cases}$$

Sectors defined by:

$$\Theta_{\mathrm{sct},j} = \theta(\min\{Q_{\mathrm{res},i}^2\} - Q_{\mathrm{res},j}^2)$$