

PHY 117 HS2024

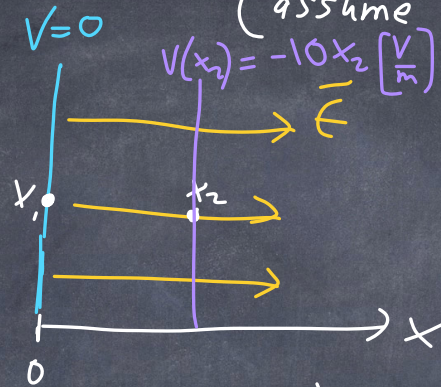
Today: Electric potential
Capacitance
Dielectrics
Electrodynamics { Electric current
{ Resistance

Note: sometimes we refer to
a potential difference ΔV as V

Be careful with E_0 + ϵ_0 :
(two different things)

Week 9, Lecture 2
Nov. 13th, 2024
Prof. Ben Kilminster

what is $V(x)$ if $\vec{E} = 10 \frac{N}{C} \hat{x}$? $E = 10 \frac{N}{C} = 10 \frac{V}{m}$
 (assume that $V=0$ at $x=0$)



$$dV = -\vec{E} \cdot d\vec{l}$$

$$dV = -10 \frac{V}{m} \hat{x} \cdot dx \hat{x}$$

$$dV = -10 dx \left[\frac{V}{m} \right]$$

$$V(x_2) - V(x_1) = \int_{x_1}^{x_2} dV = \int_{x_1}^{x_2} -10 dx \left[\frac{V}{m} \right] = -10x \left[\frac{V}{m} \right]_{x_1}^{x_2}$$

$$= 10 \frac{V}{m} (x_1 - x_2)$$

we are told that $V=0$ at $x=0$

$$\text{so } V(x_2) - V(x_1) = 10 \frac{V}{m} (x_1 - x_2)$$

$$\text{since } V(x_1=0) = 0$$

$$\Rightarrow V(x_2) - 0 = 10 \frac{V}{m} (0 - x_2) \Rightarrow V(x_2) = -10x_2 \left[\frac{V}{m} \right]$$

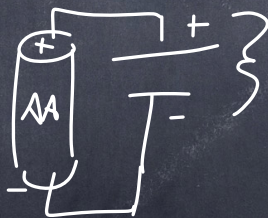
we saw that $\Delta U = q \Delta V$

\uparrow energy [J] \uparrow charge [C] \uparrow potential [V]

A convenient unit of energy is the electron volt
 example $[e \cdot V]$

$$\Delta U = \underset{\substack{\uparrow \\ \text{charge} \\ \text{of electron}}}{1 e} V = 1 (1.6 \times 10^{-19} \text{ C})(1 \text{ V}) = 1.6 \times 10^{-19} \text{ J}$$

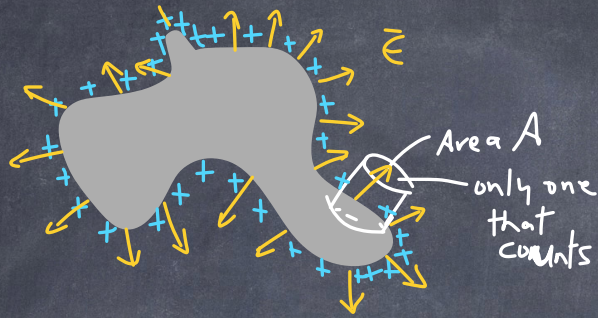
battery



$\Delta V = 1.5 \text{ V} \Rightarrow$ An electron moving through 1.5 V of potential difference will gain 1.5 eV of energy

The Large Hadron Collider (LHC) at CERN has protons with energy $6.5 \text{ TeV} = 6.5 \times 10^{12} \text{ eV}$

What is \vec{E} close to a conductor?



Near surface, $\vec{E} \parallel \hat{n}$
Also, for a conductor, charge is on surface,

$$\oint_S \vec{E} \cdot \hat{n} dA = \frac{Q}{\epsilon_0}$$

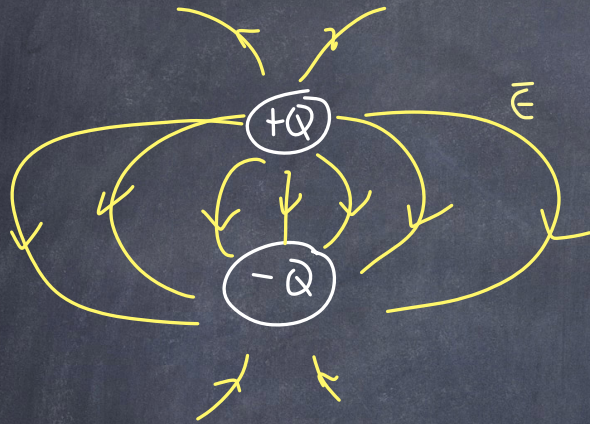
$$\sigma = \frac{Q}{A} \quad Q = \sigma A$$

$$E A = \frac{\sigma A}{\epsilon_0}$$

$$\vec{E} = \frac{\sigma}{\epsilon_0} \hat{n} \quad \text{normal to surface}$$

Capacitors :

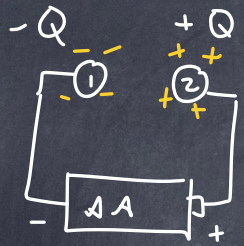
Any two oppositely charged conductors form a capacitor



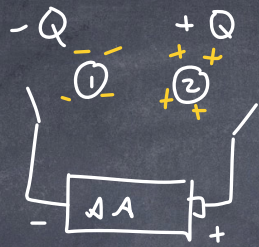
In a circuit diagram, the symbol for a capacitor is $\text{---}||\text{---}$ or $\text{---}|\text{---}$

we can charge a capacitor by connecting each conductor to the opposite terminals of a battery.

Let's charge a capacitor:



we charge with
 $\Delta V = 1.5 \text{ V}$



Q on 1 is opposite to Q on 2.
There is a potential difference between
the conductors, ΔV

we define the capacitance as $\frac{Q}{|\Delta V|} \equiv C$

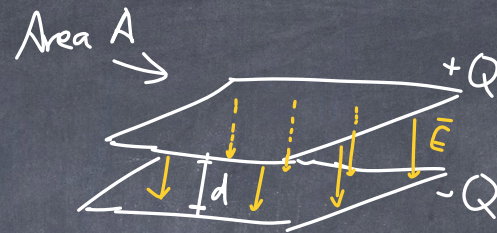
The capacitance is the ability of 2 conductors
to store electrical energy.

The units of capacitance are
1 Farad = 1 F = $1 \frac{\text{C}}{\text{V}}$

Energy stored in a capacitor:

$$U = \frac{1}{2} C (\Delta V)^2$$

what is the capacitance of a parallel-plate capacitor, with area A , separation d , and a charge $+Q$, $-Q$



we know that $E = \frac{\sigma}{\epsilon_0}$, $\sigma = \frac{Q}{A} \Rightarrow E = \frac{Q}{\epsilon_0 A}$ \oplus

The field is uniform between the 2 plates.

$$|\Delta V| = \left| -\int_0^d \vec{E} \cdot d\vec{l} \right| = E d = \frac{Q d}{\epsilon_0 A}$$

$$C = \frac{Q}{|\Delta V|} = \frac{\cancel{Q}}{\frac{\cancel{Q} d}{\epsilon_0 A}} = \frac{\epsilon_0 A}{d}$$

$$C = \frac{\epsilon_0 A}{d}$$

capacitance only depends on the dimensions.



← capacitance is $C = \frac{\epsilon_0 A}{d} = \frac{Q}{|\Delta V|}$

$$|\Delta V| = \frac{Qd}{\epsilon_0 A}$$

If we put Q on our capacitor, and increase d , then ΔV increase

Energy stored in a capacitor

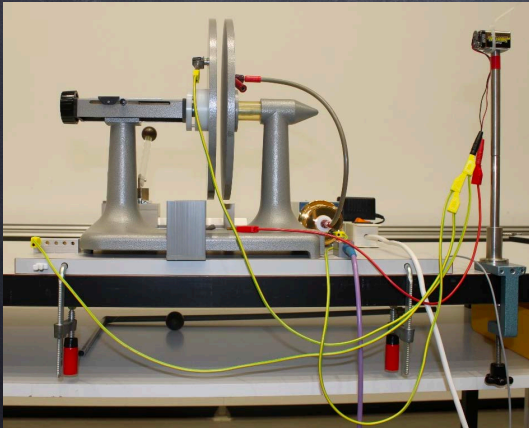
$$U = \frac{1}{2} C (\Delta V)^2$$

$$U = \frac{1}{2} \frac{\epsilon_0 A}{d} E^2 d^2$$

$$U = \frac{1}{2} \epsilon_0 A E^2 d$$

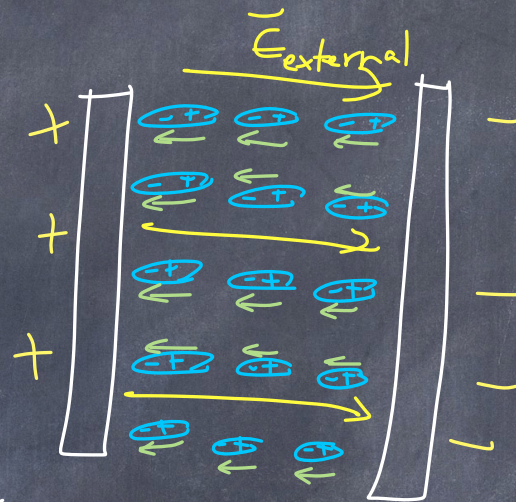
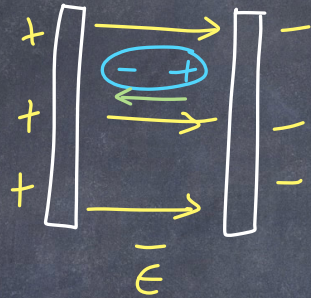
$$\Delta U = \frac{1}{2} \epsilon_0 A E^2 (\Delta d) = -W$$

we do work to increase the stored energy of the capacitor



Dielectrics - In a non-conducting material, we know that dipoles tend to align opposite to the \vec{E} -field

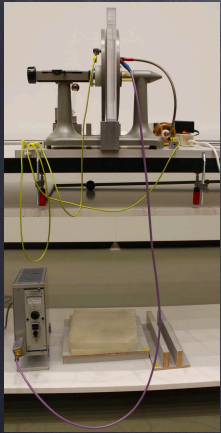
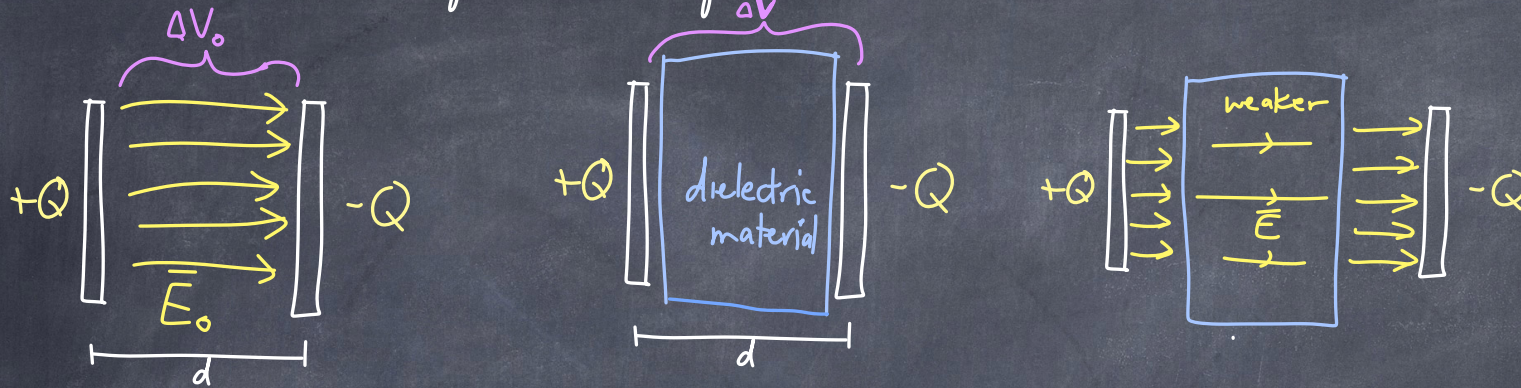
\vec{E} -field of the dipole itself



The \vec{E} -fields of the dipoles sum up

$$\text{total } \vec{E} \text{ inside} = \vec{E}_{\text{external}} - \vec{E}_{\text{dipoles}}$$

Parallel-plate capacitor with a dielectric



Initially, we have $\Delta V_0 + \epsilon_0$ with no dielectric
 $|\Delta V_0| = E_0 d$

And we know that $E = \frac{E_0}{k_d}$, so

we can get $\Delta V = E d = \frac{E_0 d}{k_d} = \frac{|\Delta V_0|}{k_d} \Rightarrow$ Adding dielectric decrease ΔV

So
$$C = \frac{Q}{|\Delta V|} = \frac{Q k_d}{V_0} = k_d C_0$$

for our parallel plate capacitor:

$$C_0 = \frac{\epsilon_0 A}{d} \Rightarrow C = k_d \frac{\epsilon_0 A}{d}$$

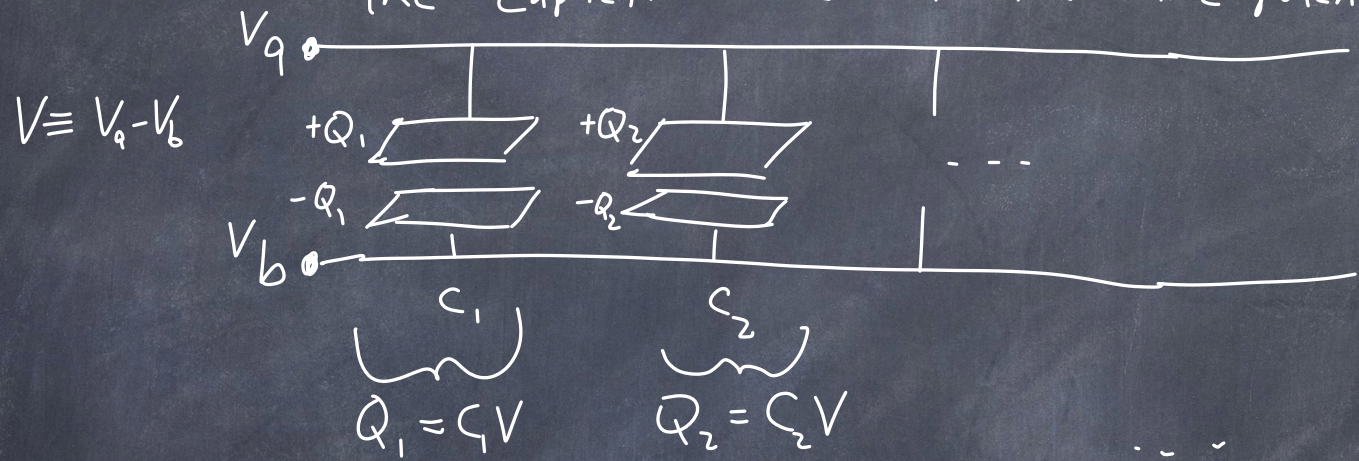
In general:

$$C = Q/V$$

$$Q = CV$$

Combining capacitors:

when we combine capacitors in parallel,
the capacitors are at the same potential



The total charge stored is $Q = Q_1 + Q_2 = C_1 V + C_2 V$

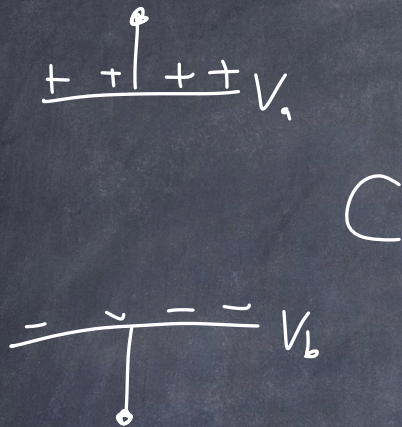
so $Q = (C_1 + C_2) V$

The equivalent capacitance of parallel capacitors

is

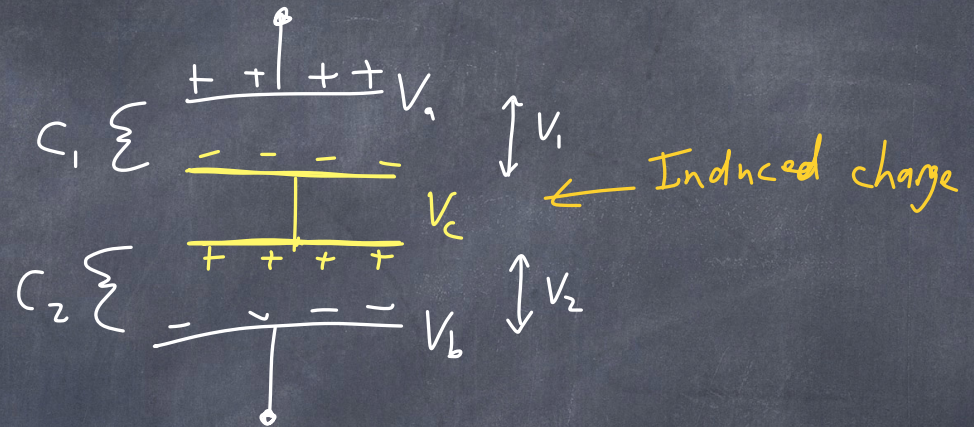
$$C_{eq} = C_1 + C_2 + \dots = \frac{Q}{V}$$

Combining capacitors in series:



$$V \equiv V_a - V_b$$

$$V = \frac{Q}{C}$$



$$V_1 = V_a - V_c \quad V_2 = V_c - V_b$$

$$V_1 = \frac{Q}{C_1} \quad V_2 = \frac{Q}{C_2}$$

The total potential difference is $V = V_a - V_b$

$$\text{so } V = V_a - V_b = (V_a - V_c) + (V_c - V_b) = \frac{Q}{C_1} + \frac{Q}{C_2}$$

$$V = Q \left(\frac{1}{C_1} + \frac{1}{C_2} + \dots \right)$$

In general,

$$\frac{1}{C_{eq}} = \frac{1}{C_1} + \frac{1}{C_2} + \dots$$

rule for adding capacitors
in series.

$$\text{If } C_1 = 2\mu\text{F} + C_2 = 2\mu\text{F}$$

stronger \rightarrow In parallel: $C_{eq} = C_1 + C_2 = 4\mu\text{F}$

weaker \rightarrow In series: $\frac{1}{C_{eq}} = \frac{1}{2\mu\text{F}} + \frac{1}{2\mu\text{F}} \Rightarrow C_{eq} = 1\mu\text{F}$

Since $U = \frac{1}{2} C \Delta V^2 \Rightarrow$ stored energy is more when capacitors are in parallel



Capacitors in parallel
vs.
in series

↑
more energy,
more spark.

Electrostatics: $E=0$ in a conductor, V is the same everywhere in a conductor.
(static charges)

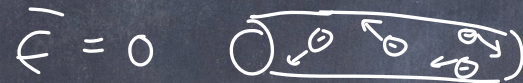
Electrodynamics: (moving charges) Current: rate of flow of electric charge.

current $\rightarrow I = \frac{\Delta Q}{\Delta t}$
 \leftarrow charge
 \leftarrow time

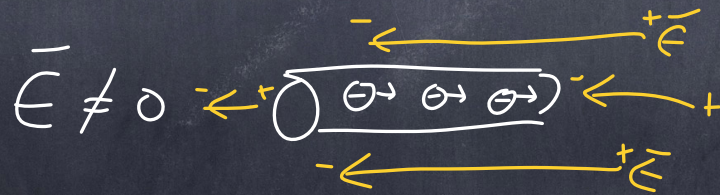
units $1\left[\frac{C}{s}\right] = 1A$

Amp = Ampere = A

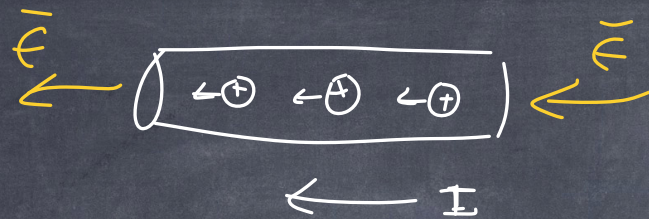
consider a conductor



Free electrons in conductor move randomly due to thermal movement
 $v \sim 10^6$ m/s



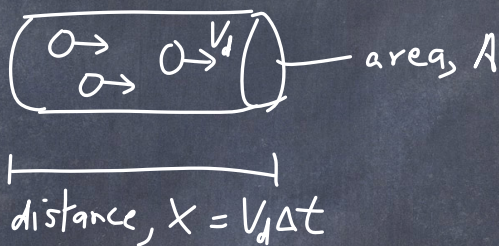
Electrons move in direction opposite to the \vec{E} -field



By convention, the current I moves in the direction that positive charges would move. (I points in same direction as \vec{E})

Calculate speed of charge carriers:

$$n = \frac{\# \text{ charge carriers}}{\text{volume}}$$



v_d : drift velocity
 volume = $A \cdot x = A v_d \Delta t$

Total charge in this cylinder is:

$$\begin{aligned} \Delta Q &= \frac{\# \text{ charge carriers}}{\text{volume}} \cdot \text{volume} \cdot \text{charge per carrier} \\ &= n \cdot \text{volume} \cdot e \end{aligned}$$

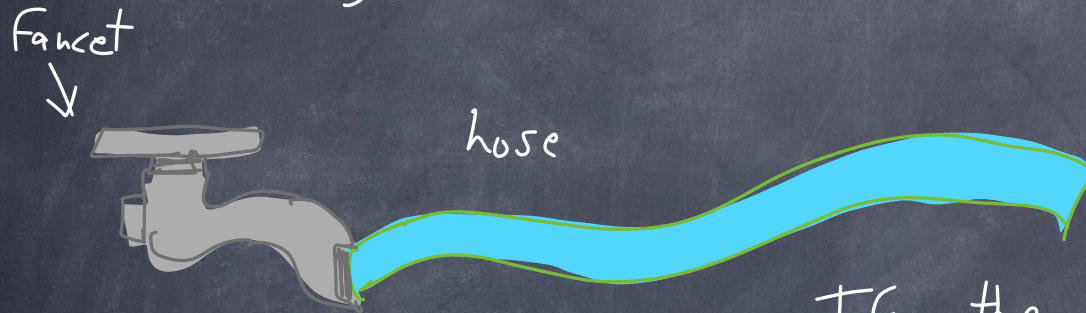
$$\Delta Q = n A v_d \Delta t e$$

$$\text{so } I = \frac{\Delta Q}{\Delta t} = \frac{n A v_d \cancel{\Delta t} e}{\cancel{\Delta t}}$$

$$v_d = \frac{I}{n A e} \sim 10^{-5} \frac{\text{m}}{\text{s}}$$

If v_d is so slow, why does electricity seem so fast?

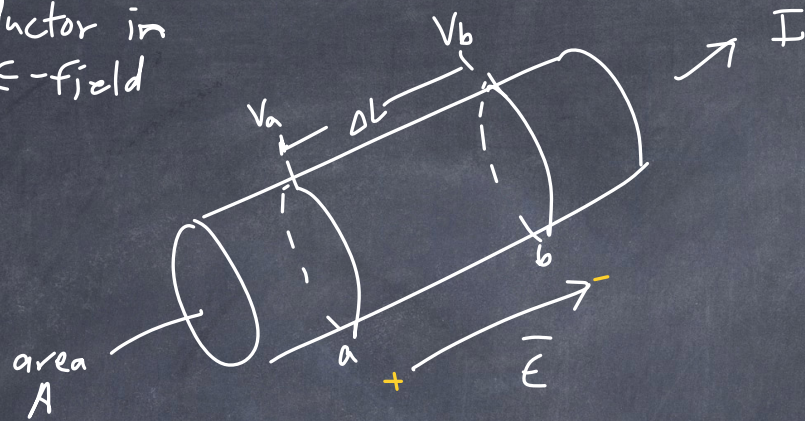
Analogy to a hose with water in it.



If the hose is already filled with water, then when you turn it on, since the fluid is incompressible, the flow is instantaneous.

Electricity is \sim instantaneous despite the slow electrons.

Conductor in
an E -field



we know that

$$V = V_a - V_b = E \Delta L$$

$$V_a - V_b > 0$$

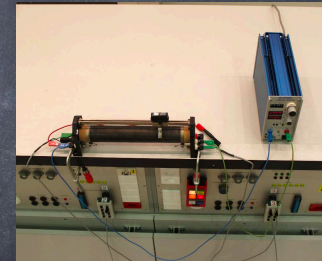
V is the potential difference

For most materials, the current is
proportional to the potential difference:

$$I \propto V$$

This constant of proportionality is $\frac{1}{R}$
where R is the resistance of the material

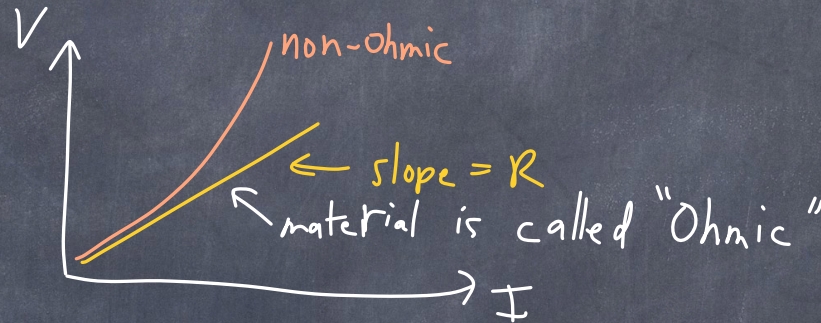
$$I = \left(\frac{1}{R}\right)V \Rightarrow V = IR$$



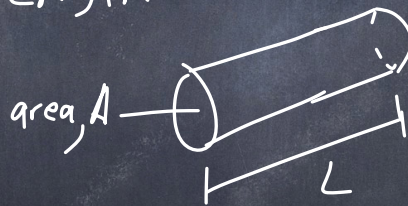
R has units of Ω (ohms)

where $1 \Omega = 1 \frac{V}{A} = \frac{1 \text{ volt}}{\text{Amp}}$

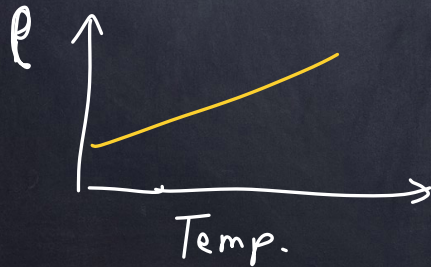
$V = IR$



The resistance R depends on the material, area, length



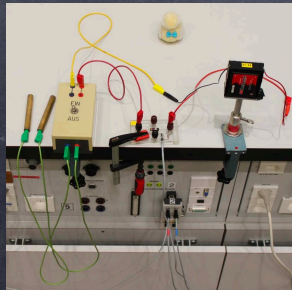
$R = \frac{L}{A} \rho$ " ρ " is the resistivity
per material



$\rho = \rho_{20} \left[1 + \alpha (t_c - 20^\circ) \right]$
resistivity at 20°C α per material t_c Temp. in Celsius

<u>material</u>	<u>$\rho_{20} [\Omega \cdot m]$</u>	<u>$\alpha \left[\frac{1}{\text{C}} \right]$</u>
Copper	$1.7 \text{ E-}8$	$3.9 \text{ E-}3$
aluminum	$2.8 \text{ E-}8$	$3.9 \text{ E-}3$
wood	$10^8 - 10^{14}$	
glass	$10^{10} - 10^{14}$	

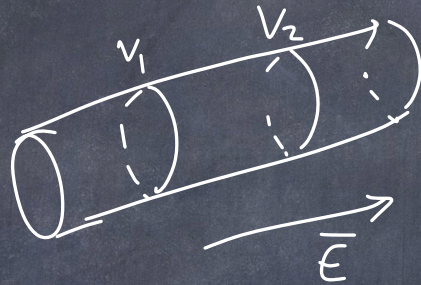
IF R is small, then $I = \frac{V}{R}$ is big:
 more current in a copper tube
 than in a wooden one.



Sometimes people use

$$\text{Conductivity} = \frac{1}{\text{resistivity}}$$

Energy is lost in a conductor as electrical energy is converted into thermal energy.



Potential decrease from V_1 to V_2
 $V = V_1 - V_2$

Loss in energy $\Delta U = \Delta Q (V_2 - V_1)$

ΔQ is the amount of charge flowing from V_1 to V_2

In time Δt , $\frac{\Delta U}{\Delta t} = \frac{\Delta Q(V)}{\Delta t}$

↑
power
[Watts]

↑
I
[Amps]

↑
[Volts]

V is the potential difference here

1 Watt = 1 Amp · Volt = 1 A · V

$$P = IV$$

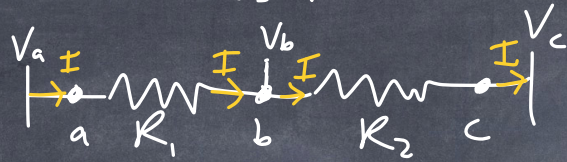
power

Since we know $V = IR$ ↑

$$P = I^2 R = \frac{V^2}{R}$$

The energy loss depends on R .
 At constant I , a higher R produces more heat generated.
 At constant V , a higher R produces less heat.

Resistors in series:



what goes in, must come out
 $I_a = I_b = I_c = I$

Equivalent resistance

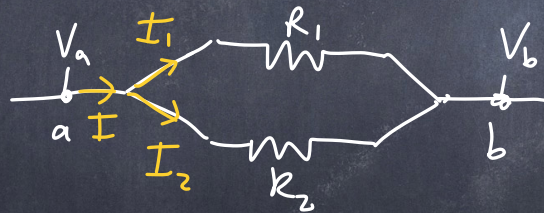
$$R_{eq} = R_1 + R_2 + \dots$$

$$V_b = V_a - IR_1$$

$$V_c = V_a - IR_1 - IR_2$$

potential decreases,
but
the current
stays the same.

Resistors in parallel:



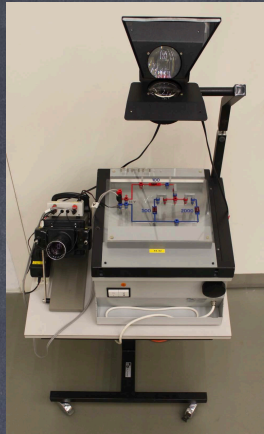
$$I = I_1 + I_2$$

Equivalent resistance
decreases (more ways for
the current to flow)

$$\frac{1}{R_{eq}} = \frac{1}{R_1} + \frac{1}{R_2} + \dots$$

voltage drop $V_a - V_b$ is the
same across both pathways

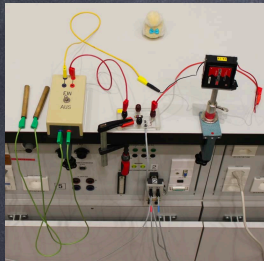
$$V_{ab} = V_a - V_b = I_1 R_1 = I_2 R_2$$



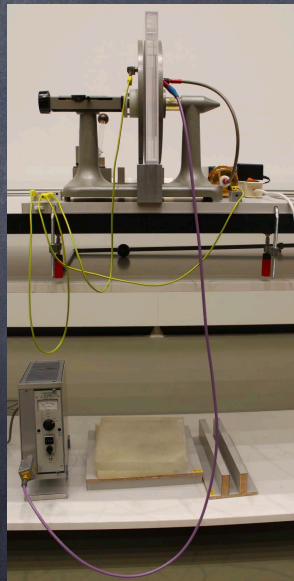
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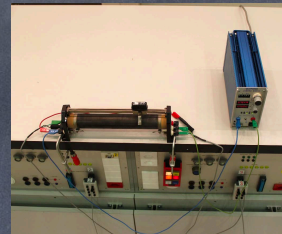
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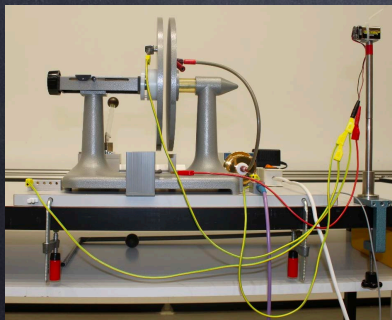
ES70



ES44



ES61



ES34