5 Precision measurements in pion decay

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Within the Standard Model the decay $\pi \to e \nu$ is helicity suppressed $(\Gamma_{\pi \to e \nu}/\Gamma_{\pi \to \mu \nu}) \approx m_e^2/m_\mu^2/(1-m_\mu^2/m_\pi^2)^2 \approx 1.2 \times 10^{-4})$ which makes it sensitive to a number of hypothetical exotic contributions. In the radiative decay $\pi \to e \nu \gamma$ the trivial inner-bremsstrahlung contribution is helicity suppressed as well and structure-dependent contributions dominate in large parts of the observable phase space. For the decay $\pi \to e \nu$ data taking just started. New results from our $\pi \to e \nu \gamma$ measurements have recently become available (1).

5.1 Testing lepton universality, the $\pi^+ \rightarrow e^+ \nu$ decay

The experiment aims at improving the accuracy in the determination of the $\pi \to e \nu/\pi \to \mu \nu$ branching ratio by an order of magnitude (2). See past year's annual report for an introduction to the concept of lepton universality and a discussion of theoretical scenarios that might result in violations.

During a three-months measuring period in 2007 feasibility tests were performed at a pion stop rate of only $10^3/s$. Our institute supplied a new set of beam counters and the associated 8 GHz transient recorders. We also refurbished the cylindrical scintillator hodosope situated in between the tracking detectors and the spherical Csl calorimeter. A preliminary analysis revealed no show-stoppers and the PSI program advisory committee gave green light to a first production run in 2008. By the time of writing data are taken at a $5 \times 10^3/s$

stop rate which soon will be increased to $10^4/s$ giving 3×10^4 observed $\pi\to e\nu$ decays per day which should bring down the statistical error to O(10^{-3}) in 2008, significantly below the present uncertainty of the world average for the branching ratio ($1.231(4)\times10^{-4}$).

5.2 Pion substructure revealed in the $\pi^+ \rightarrow e^+ \nu \gamma$ decay

In the standard description of radiative pion decay the decay amplitude has contributions from inner bremsstrahlung (IB), structure dependent radiation (SD $^\pm$), and interference terms. The structure dependent amplitude is parameterized by F_V and F_A , vector and axial vector form factors, respectively (see Ref. (3) for a comprehensive review). The conserved vector current (CVC) hypothesis relates F_V to the π^0 lifetime, yielding F_V =0.0259(9). Chiral symmetry (ChPT) calculations (4; 5; 6) give a value for F_A in the range 0.010–0.014.

Early experimental studies were limited by relatively low event statistics. The authors evaluated $\gamma \equiv F_A/F_V$ resulting in overall uncertainties ranging from 12% to 56% (7; 8; 9; 10; 11; 12). None of these measurements considered the dependence of the form factors on the invariant mass of the $\mathrm{e}^+\nu$ pair.

During the years 1999 - 2001 a first data set with a 16 fold pre-scaled trigger was collected as a byproduct of the PIBETA $\pi^+ \to \pi^0 \mathrm{e}^+ \nu$ branching ratio measurement (13). The analy-

sis of the data resulted in a significant improvement in the accuracy of γ =0.443(15) (14) but also noted a substantial deficit of observed events in the high- E_{γ} /low- $E_{\rm e^{+}}$ kinematic region. For this reason a dedicated experiment was performed during 2004 after Zürich joined the new PEN collaboration. During these measurements the pion stop rate was reduced by almost an order of magnitude to $\sim 10^{5}$ /s. The reduced stop rate resulted in a dramatic reduction of accidental $e\gamma$ coincidences which allowed us to observe the decays in a much larger region of phase space.

In order to keep accidental background under control the trigger for data readout required at least one cluster of Csl crystals with a total energy above 52 MeV, the endpoint energy in muon decay. Candidate events exhibit one neutral shower in the calorimeter coincident in time with a positron track. Three kinematic regions were investigated:

region	E_{e^+} (MeV)	E_{γ} (MeV)	events
	52 - 70	52 - 70	36k
II	10 - 52	52 - 70	16k
Ш	52 - 70	10 - 52	13k

with relative angle $\theta_{e\gamma}$ >40° for all regions. Under these conditions the decay is completely dominated by the SD⁺ structure-dependent term proportional to $(F_V + F_A)^2$ (3).

The values of F_V , F_A , and a (defined by $F_V(q^2)=F_V(0)(1+a\cdot q^2)$, $F_A(q^2)=F_A(0)$, and $q^2=1-(2E_\gamma/m_\pi)$) were fit to the measured data by χ^2 minimization:

$$F_A = +0.0117(17)$$

 $F_V = +0.0258(17)$
 $a = +0.095(58)$

which represent fourteen-fold a improvement over previous measurements. agreement between the The excellent measured calculated distributions and the best-fit parameter values

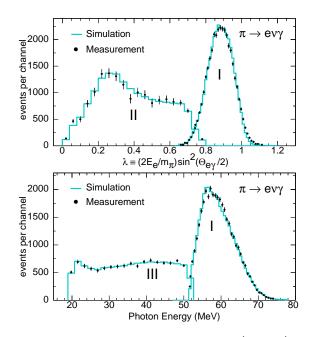


Figure 5.1: Background-subtracted $\pi^+ \to e^+ \nu \gamma$ distribution of the kinematic variable $\lambda \equiv (2 E_e/m_\pi) \sin^2(\theta_{e\gamma}/2)$ for regions I and II (top panel), and of the photon energy for regions I and III (bottom panel). For the simulation the best-fit values of F_V , F_A , and a were used.

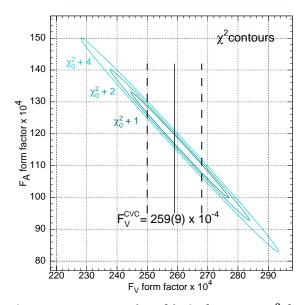


Figure 5.2: Contour plot of loci of constant χ^2 for the minimum value χ^2_0 plus 1, 2, and 4 units, respectively, in the $F_V \times F_A$ parameter plane, keeping the parameter a=0.041. The range of the CVC prediction, $F_V = 0.0259(9)$ is indicated.

Fig. 5.1) reflects the perfect description of the data by CVC and standard V-A theory in all regions of phase space. As discussed above the data are most sensitive to the combination $F_A + F_V$ which is illustrated in Fig. 5.2.

Conclusions

- The best-fit value of F_A agrees well with ChPT calculations, tending to the middle of the reported range (4; 5; 6).
- A more precise measurement of $\tau_{\pi 0}$ is needed to make full use of our data in determining F_A .
- The deficit of observed events in the high- $E_{\gamma}/{\rm low}\text{-}E_{\rm e^+}$ has not been confirmed.
- The branching ratio for E $_{\gamma}$ > 10 MeV and $\theta_{e\gamma}$ > 40° is measured to be 73.9(5)×10⁻⁸.
- The pion polarizability can be determined as follows: using the one-parameter fit we obtain $\alpha_E = -\beta_M = (2.78 \pm 0.02_{exp} \pm 0.10_{Fv}) \times 10^{-4} \, \mathrm{fm^3}$, where the first uncertainty comes from the fit and the second from the current CVC-derived value of F_V ; alternatively, one obtains $2.7^{+0.6}_{-0.5} \times 10^{-4} \, \mathrm{fm^3}$ based on our unconstrained fit of F_A and F_V .
- In addition, we use the latter fit result and CVC to make an independent determination of the neutral pion lifetime: $\tau_{\pi0}$ =(8.5 \pm 1.1) \times 10⁻¹⁷ s.

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